

Superconducting properties of tin-based ENZ and hyperbolic metamaterials

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Recent experiments have demonstrated that the superconducting critical temperature may be improved in various metamaterial superconductor geometries. Here, we present the results of a study of tin-based metamaterial superconductors in the epsilon-near-zero (ENZ) and hyperbolic metamaterial configurations. It was observed that T_c enhancement is significantly reduced when the metamaterial structural dimensions exceed 240nm, the superconducting coherence length in pure tin.

1. Introduction

Designing new materials with enhanced superconducting properties remains one of the most important goals of condensed matter physics. A strategy that deliberately engineers the dielectric properties of nanostructured composite superconductor-dielectric metamaterials has been demonstrated to result in an enhanced electron pairing interaction that increases the value of the critical temperature, T_c , of so-called “metamaterial superconductors” [1-5]. For example, Al_2O_3 -coated aluminum nanoparticles can be used to form epsilon near zero (ENZ) core-shell metamaterial superconductors with a T_c that is three times that of pure aluminum [4]. Alternatively, a multilayer $\text{Al}/\text{Al}_2\text{O}_3$ hyperbolic metamaterial structure may be fabricated [5], which exhibits similar T_c enhancement, while having much better transport properties. Hyperbolic metamaterials are extremely anisotropic uniaxial materials, which behave like a metal ($\epsilon_{xx} = \epsilon_{yy} < 0$) in one direction and like a dielectric ($\epsilon_{zz} > 0$) in the orthogonal direction [6]. The hyperbolic metamaterial geometry shown in Fig. 1 is a structure based on periodic metallic layers separated by dielectric layers. As was originally noted in [7], typical high T_c superconductors (such as BSCCO) exhibit hyperbolic behavior in the far infrared and THz frequency ranges, which may be partially responsible for their very large T_c values.

As has been noted above, most of the initial metamaterial superconductor geometries were successfully demonstrated using aluminum as a superconducting component of the composite metamaterial. Aluminum is ideal for these proof of principle experiments because its critical temperature is quite low ($T_{c \text{ Al}} = 1.2 \text{ K}$ [8]), leading to a very large superconducting coherence length $\xi = 1600 \text{ nm}$ [8]. Such a large

value of ξ facilitates the metamaterial fabrication requirements since the validity of macroscopic electrodynamic description of the metamaterial superconductor requires that its spatial structural parameters must be much smaller than ξ . In this paper we will expand this approach to another elemental superconductor, tin, which also has rather large superconducting coherence length $\xi = 240$ nm [8]. We have realized both ENZ and hyperbolic metamaterial geometries based on tin, and observed the effect of T_c increase in the tin-based metamaterial superconductors. In addition, we studied the effect of metamaterial component dimensions on T_c in the tin-based ENZ metamaterials demonstrating that T_c enhancement is significantly reduced when the metamaterial structural parameter exceeds 240 nm, the coherence length of pure tin. This observation supports the validity of the macroscopic electrodynamic (metamaterial) approach to designing new materials with enhanced superconducting properties.

2. Macroscopic electrodynamic description of superconductivity

Electromagnetic properties play a very important role in the electron pairing in superconductors [9]. According to the BCS theory, the Cooper pairing mechanism may be described as an attractive interaction of electrons which results from the polarization of the ionic lattice created by motion of these electrons through the lattice. Based on this physical picture, Kirzhnits *et al.* formulated their description of superconductivity in terms of the dielectric response function of the superconductor [9]. Within the scope of this “macroscopic electrodynamics” approach, the electron-electron interaction in a superconductor can be expressed in the form of an effective Coulomb potential

$$V(\vec{q}, \omega) = \frac{4\pi e^2}{q^2 \varepsilon_{\text{eff}}(\vec{q}, \omega)} = V_C \frac{1}{\varepsilon_{\text{eff}}(\vec{q}, \omega)} \quad (1)$$

where V_C is the Fourier-transformed Coulomb potential in vacuum, and $\varepsilon_{\text{eff}}(q, \omega)$ is the dielectric response function of the superconductor, which is treated as an effective medium. Following this formalism, it appears natural to use the recently developed electromagnetic metamaterial toolbox [10] to engineer the $\varepsilon_{\text{eff}}(q, \omega)$ of an artificial “metamaterial superconductor”, so that the electron pairing interaction described by Eq.(1) will be maximized. Indeed, it was predicted in [1, 2] that a considerable enhancement of the electron pairing interaction may be expected in the ENZ [11] and hyperbolic metamaterial scenarios [6]. In both cases $\varepsilon_{\text{eff}}(q, \omega)$ may become small and negative in substantial portions of the four-momentum (q, ω) space near the electron Fermi surface, leading to considerable enhancement of the electron pairing interaction. For example, in the case of a hyperbolic metamaterial the effective Coulomb potential from Eq. (1) assumes the form

$$V(\vec{q}, \omega) = \frac{4\pi e^2}{q_z^2 \varepsilon_2(\vec{q}, \omega) + (q_x^2 + q_y^2) \varepsilon_1(\vec{q}, \omega)} \quad (2)$$

where $\varepsilon_{xx} = \varepsilon_{yy} = \varepsilon_1$ and $\varepsilon_{zz} = \varepsilon_2$ have opposite signs. As a result, the effective Coulomb interaction of two electrons may become attractive and very strong along spatial directions where

$$q_z^2 \varepsilon_2(\vec{q}, \omega) + (q_x^2 + q_y^2) \varepsilon_1(\vec{q}, \omega) \approx 0 \quad (3)$$

As summarized in [12], the critical temperature of a superconductor in the weak coupling limit is typically calculated as

$$T_c = \theta \exp\left(-\frac{1}{\lambda_{\text{eff}}}\right), \quad (4)$$

where θ is the characteristic temperature for a bosonic mode mediating electron pairing (such as the Debye temperature θ_D in the standard BCS theory), and λ_{eff} is the dimensionless coupling constant defined by $V(q, \omega) = V_C(q) \varepsilon^{-1}(q, \omega)$ and the density of states ν (see for example [13]):

$$\lambda_{\text{eff}} = -\frac{2}{\pi} \nu \int_0^\infty \frac{d\omega}{\omega} \langle V_C(q) \text{Im} \varepsilon^{-1}(\vec{q}, \omega) \rangle, \quad (5)$$

where V_C is the unscreened Coulomb repulsion, and the angle brackets denote average over the Fermi surface. Additional zeroes of the effective dielectric response function $\varepsilon_{\text{eff}}(q, \omega)$ of the metal-dielectric metamaterial described by Eq. (3) lead to increase of λ_{eff} compared to the bulk metal, resulting in higher T_c of the metamaterial. Very recently it was demonstrated that a multilayer Al/Al₂O₃ hyperbolic metamaterial structure indeed exhibits considerable T_c enhancement compared to the critical temperature in bulk aluminum [5]. Here, we report that similar ENZ and hyperbolic metamaterial configurations increase T_c in tin-based metamaterial superconductors.

3. Tin-based ENZ superconductors: metamaterial structural parameter size effect

According to the Maxwell-Garnett approximation [14], mixing of nanoparticles of a superconducting “matrix” with dielectric “inclusions” (described by the dielectric constants ε_m and ε_d , respectively) results in the effective medium with a dielectric constant ε_{eff} , which may be obtained as

$$\left(\frac{\epsilon_{eff} - \epsilon_m}{\epsilon_{eff} + 2\epsilon_m} \right) = (1-n) \left(\frac{\epsilon_d - \epsilon_m}{\epsilon_d + 2\epsilon_m} \right), \quad (6)$$

where n is the volume fraction of metal. An explicit expression for the inverse dielectric response function ϵ_{eff}^{-1} of this effective medium may be written as

$$\epsilon_{eff}^{-1} = \frac{n}{(3-2n)} \frac{1}{\epsilon_m} + \frac{9(1-n)}{2n(3-2n)} \frac{1}{(\epsilon_m + (3-2n)\epsilon_d/2n)}, \quad (7)$$

which makes it apparent that ϵ_{eff}^{-1} of the mixture acquires an additional pole at

$$\epsilon_m \approx -\frac{3-2n}{2n} \epsilon_d, \quad (8)$$

(in addition to the bulk metal plasmon pole at $\epsilon_m \approx 0$). This additional pole, which has been recently identified as a plasmon-phonon pole of the metal-dielectric mixture [15], contributes to the increased coupling constant λ_{eff} of the metamaterial, which leads to increased T_c .

The ENZ behavior described above, and the corresponding increase of the superconducting critical temperature has been recently observed in Sn/BaTiO₃ nanoparticle mixtures [3]. The metamaterial samples in these experiments were prepared using commercially available tin and barium titanate nanoparticles obtained from the US Research Nanomaterials, Inc. The average diameter of the BaTiO₃ nanoparticles was 50 nm, while tin nanoparticle size was specified as 60-80 nm. Since both nanoparticle sizes fall substantially below the superconducting coherence length in pure tin $\xi_{Sn} \sim 230$ nm [8], the metamaterial approach seems to be well justified in this case. The Sn/BaTiO₃ superconducting metamaterials were fabricated by combining the given amounts of Sn and BaTiO₃ nanoparticle powders by volume into a single test tube filled with de-ionized water. The resulting suspensions were sonicated and

magnetically stirred for 30 minutes, followed by water evaporation. The remaining mixtures of Sn and BaTiO₃ nanoparticles were compressed into test pellets using a hydraulic press. The ENZ conditions in these initial experiments were achieved in the 30-50% range of the volume fraction of BaTiO₃. The T_c increased from the pure Sn value of 3.68 K with increasing BaTiO₃ concentration to a maximum $\Delta T_c \sim 0.15$ K (or 4%) compared to the pure tin sample for the 40% sample followed by T_c decreasing at higher volume fractions [3].

The commercial availability of a broad variety of tin and BaTiO₃ nanoparticle sizes (in the 50 nm to several micrometers range) allowed us to further extend this study to larger nanoparticle sizes, and verify that the observed T_c increase indeed gradually diminishes with an increase in particle sizes beyond the superconducting coherence length in pure tin $\xi_{Sn} \sim 230$ nm. These experiments provide a very strong support for the metamaterial superconductor concept.

A similar metamaterial fabrication technique described above has been applied to samples containing larger nanoparticle sizes. Typical Scanning Electron Microscopy (SEM) images of the resulting Sn/BaTiO₃ composite metamaterials are shown in Fig. 2(a). The original compressed nanoparticles are clearly visible in both images. The resulting average metamaterial composition has been verified after fabrication using an SEM with energy dispersive X-ray spectroscopy (EDS). An example of such compositional analysis spectra is shown in Fig.2(b). The measured individual element abundancies were consistent with this nominal composition. Such EDX spectra were used to establish the homogeneous character of the fabricated composite metamaterials.

The ENZ character of the fabricated Sn/BaTiO₃ metamaterial superconductor samples has been verified by infrared reflectance measurements of the samples in an

extended frequency and temperature range. The real and imaginary parts of the dielectric function were obtained by Kramers-Kronig analysis of the infrared reflectance data constrained by ellipsometry data at higher frequencies [16]. As plotted in Fig. 3, the real part of the frequency dependent dielectric constant of the metamaterial samples at 5 K and at room temperature indeed appeared to be much smaller in magnitude compared to the real part of the dielectric constant of tin.

The superconducting critical temperature T_c of various Sn/BaTiO₃ metamaterial samples was measured as a function of nanoparticle sizes and volume fractions of BaTiO₃ via the onset of diamagnetism using a MPMS SQUID magnetometer. The zero field cooled (ZFC) magnetization per unit mass for several samples with varying concentrations of BaTiO₃ for two different sets of nanoparticle sizes is plotted in Fig. 4. The effect of the superconducting critical temperature T_c increase appears to be strongly suppressed in metamaterial superconductors made with particle sizes considerably larger than the superconducting coherence length in pure tin $\xi_{\text{Sn}} \sim 230$ nm.

The temperatures of the onset of the superconducting transition for samples made with various nanoparticle sizes as a function of BaTiO₃ volume fraction are plotted in Fig. 5(a). The pure tin samples were prepared from pressed tin nanoparticles of the same diameter using the same method of preparation. Note that the value of T_c for these samples agreed with the expected value of pure Sn. The increase of T_c for the samples with 40% volume fraction of BaTiO₃ measured as a function of nanoparticle size is plotted in Fig. 5(b). Once again, T_c increase appears to be strongly suppressed in metamaterial superconductors made with particle sizes considerably larger than the superconducting coherence length in pure tin. These results validate the metamaterial approach to dielectric response engineering in metamaterial superconductors.

4. Tin-based hyperbolic metamaterial superconductors

As has been mentioned above, the hyperbolic metamaterial strategy also appears to be a very effective tool of dielectric response engineering in metamaterial superconductors. We were able to validate this approach with tin-based hyperbolic metamaterials. In the hyperbolic metamaterial geometry shown schematically in Fig.1 the role of metal layers was played by thin tin films, while either CaF_2 or MgF_2 was used as the dielectric layer material. The individual layers of Sn and CaF_2 were fabricated sequentially using thermal evaporation. This process was repeated to create multilayer Sn- CaF_2 hyperbolic metamaterial samples on glass substrates. For the metamaterial approach to be successful, the thicknesses of superconductor and dielectric layers should be much smaller than the superconducting coherence length in pure tin $\xi_{\text{Sn}} \sim 230$ nm, and the dielectric layer thickness must be thin enough to allow electron tunneling between the tin layers. The increase of T_c in hyperbolic metamaterials also depends on the volume fraction of metal in metal/dielectric multilayers [5]. Therefore, we tried multiple layer deposition thicknesses to find the optimal tin/ CaF_2 ratio in the metamaterial. The actual thicknesses of the tin and CaF_2 films in the nanofabricated samples were then determined by Atomic Force Microscopy (AFM).

To demonstrate that our multilayer samples exhibit hyperbolic behaviour, we studied their optical properties using polarization reflectometry (Fig. 6). The theoretical values of the ϵ_1 and ϵ_2 components of the Sn/ CaF_2 multilayer films may be obtained using the Maxwell-Garnett approximation as follows:

$$\varepsilon_1 = n\varepsilon_m + (1-n)\varepsilon_d \quad (9)$$

$$\varepsilon_2 = \frac{\varepsilon_m \varepsilon_d}{(1-n)\varepsilon_m + n\varepsilon_d} \quad (10)$$

where n is the volume fraction of metal, and ε_m and ε_d are the dielectric permittivities of the metal and dielectric, respectively [17]. The signs of the real parts of ε_1 and ε_2 (and therefore the hyperbolic character of our samples) were determined by polarization reflectometry. The metamaterial parameters were extracted from the polarization reflectometry data as described in detail in [18]. Reflectivity for s-polarization is given in terms of the incident angle θ by

$$R_s = \left| \frac{\sin(\theta - \theta_{ts})}{\sin(\theta + \theta_{ts})} \right|^2, \quad (11)$$

where

$$\theta_{ts} = \arcsin\left(\frac{\sin \theta}{\sqrt{\varepsilon_1}}\right). \quad (12)$$

Reflectivity for p-polarization is given as

$$R_p = \left| \frac{\varepsilon_1 \tan \theta_{tp} - \tan \theta}{\varepsilon_1 \tan \theta_{tp} + \tan \theta} \right|^2, \quad (13)$$

where

$$\theta_{tp} = \arctan \sqrt{\frac{\varepsilon_2 \sin^2 \theta}{\varepsilon_1 \varepsilon_2 - \varepsilon_1 \sin^2 \theta}}. \quad (14)$$

We measured p- and s- polarized absolute reflectance of the metamaterial samples at three different laser light frequencies (488 nm, 633 nm and 1550 nm.) as indicated in

Fig. 6. The absolute reflectance values were obtained by normalization to the measured reflectance of a 150 nm gold film. The dashed lines in Fig. 6 are fits using Eqs. (11-14). ε_1 and ε_2 values obtained from the fits confirm hyperbolic character of the 8 layer Sn/CaF₂ metamaterial which consists of 5.3 nm Sn layers separated by 1 nm layers of CaF₂. It is clear that the real part of the out-of-plane dielectric function is positive, while the real part of the in-plane dielectric function is negative, which confirms the dielectric nature along z-axis and metallic nature in the xy-plane, i.e. a hyperbolic metamaterial. We have also verified that the results of these measurements demonstrate good agreement with theoretical predictions obtained from Maxwell-Garnett approximation, Eqs. (9, 10), and the reported optical properties of tin [19] and CaF₂ [20]. Comparison of the calculated and measured values of the real and imaginary parts of ε_1 and ε_2 is presented in Fig.7.

The T_c of various samples (Figs. 8, 9) were determined via four-point resistivity measurements as a function of temperature using a Physical Property Measurement System (PPMS) by Quantum Design. Fig. 8 shows measured resistivity as a function of temperature for several multilayer Sn/CaF₂ metamaterials having different thicknesses of Sn layers, while the 1 nm thickness of CaF₂ layers was kept approximately constant. The Sn/CaF₂ hyperbolic metamaterial samples were found to have a higher critical temperature than pure Sn, which strongly indicates the role of hyperbolic geometry in T_c enhancement.

We have also studied changes in T_c as a function of tin volume fraction in a set of several Sn/CaF₂ and Sn/MgF₂ metamaterial samples, as shown in Fig. 9. The quantitative behaviour of T_c as a function of n may be predicted based on the hyperbolic enhancement of the electron-electron interaction (Eq. (2)) and the density of electronic

states, ν , on the Fermi surface which experience this hyperbolic enhancement [5]. Using Eqs. (9, 10), the effective Coulomb potential from Eq. (2) may be re-written as

$$V(\vec{q}, \omega) = \frac{4\pi e^2}{q^2 \left(\frac{q_z^2}{q^2} \frac{\epsilon_d \epsilon_m}{((1-n)\epsilon_m + n\epsilon_d)} + \frac{q_x^2 + q_y^2}{q^2} (n\epsilon_m + (1-n)\epsilon_d) \right)}. \quad (15)$$

Based on Eq. (15), the differential of the product νV may be written as

$$\begin{aligned} d(\nu V) &= \frac{4\pi e^2 n \sin \theta d\theta}{q^2 \left(\frac{\epsilon_d \epsilon_m}{((1-n)\epsilon_m + n\epsilon_d)} \cos^2 \theta + (n\epsilon_m + (1-n)\epsilon_d) \sin^2 \theta \right)} = \\ &= - \frac{4\pi e^2 n dx}{q^2 \left((n\epsilon_m + (1-n)\epsilon_d) - \frac{n(1-n)(\epsilon_m - \epsilon_d)^2}{(n\epsilon_d + (1-n)\epsilon_m)} x^2 \right)}, \end{aligned} \quad (16)$$

where $x = \cos \theta$, and θ varies from 0 to π . The latter expression has two poles at

$$\epsilon_m = \left(\left(1 - \frac{1}{2n(1-n)(1-x^2)} \right) \pm \sqrt{\left(1 - \frac{1}{2n(1-n)(1-x^2)} \right)^2 - 1} \right) \epsilon_d \quad (17)$$

As the volume fraction, n , of metal is varied, one of these poles remains close to $\epsilon_m = 0$, while the other is observed at larger negative values of ϵ_m . This behaviour is similar to the appearance of an additional pole in ENZ metamaterials (see Eqs. (7, 8)). Using Eqs.(4, 5), and following the numerical integration procedure described in detail in [5], a theoretical prediction of T_c as a function of n in the Sn/CaF₂ hyperbolic metamaterials may be obtained. This behaviour is plotted in Fig. 9 as a red curve. This figure demonstrates that the experimentally measured behaviour of T_c as a function of n (which is defined by the Sn layer thickness) correlates well with the theoretical prediction. As discussed in detail in [21], smaller relative increases of T_c in tin-based

metamaterial superconductors compared to similar aluminium-based metamaterial geometries may be explained by considerably larger imaginary part ϵ_m'' of the tin dielectric constant.

5. Conclusions

Our experimental results confirm recent theoretical predictions and earlier experimental observations that the superconducting critical temperature may be improved in various metamaterial superconductor geometries. Our detailed study of tin-based metamaterial superconductors in the epsilon-near-zero (ENZ) and hyperbolic metamaterial configurations demonstrated that the metamaterial superconductor approach remains valid even as the superconductor coherence length is strongly reduced compared to the very large value of $\xi = 1600$ nm in aluminum. As expected, it was observed that the T_c enhancement effect is significantly reduced when the dimensions of the metamaterial components exceed $\xi = 240$ nm superconducting coherence length in pure tin.

Our experimental observations open up numerous other possibilities for considerable T_c enhancement in other practically important simple superconductors. However, due to their much smaller coherence length metamaterial structuring of these superconductors must be performed on a much more refined scale.

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Figure Captions

Fig. 1. Schematic geometry of a “layered” hyperbolic metamaterial. In the present work tin was used as metallic layers, while either CaF_2 or MgF_2 were used as the dielectric layers.

Fig. 2. (a) SEM images of the composite Sn/BaTiO_3 metamaterials. Individual compressed nanoparticles are clearly visible in the images. The sample compositions and Sn and BaTiO_3 nanoparticle sizes are indicated above the images. (b) Verification of the average metamaterial composition using an SEM with energy dispersive X-ray spectroscopy (EDS).

Fig. 3. Confirmation of the ENZ behavior of the Sn/BaTiO_3 metamaterial superconductor samples: The real part, ε' , of the dielectric constant of the metamaterial samples at 5 K and at room temperature appears to be much smaller in magnitude compared to the real part of the dielectric constant of tin.

Fig. 4. Temperature dependence of zero field cooled magnetization per unit mass measured in a magnetic field of 10 G for several tin- BaTiO_3 metamaterial samples with varying volume fraction of BaTiO_3 (indicated in the plot) for samples made with (a) 70 nm tin and 50 nm BaTiO_3 nanoparticles (b) 1000 nm tin and 3000 nm BaTiO_3 nanoparticles.

Fig. 5. (a) The temperatures of the onset of the superconducting transition plotted as a function of volume fraction of BaTiO_3 for samples made with various nanoparticle sizes. The lines are a guide to the eye. (b) The increase of T_c for the samples with 40% volume fraction of BaTiO_3 measured as a function of nanoparticle size. The superconducting critical temperature T_c increase appears to be strongly suppressed in

metamaterial superconductors made with particle sizes considerably larger than the superconducting coherence length in pure tin $\xi_{\text{Sn}} \sim 230$ nm.

Fig. 6. Polarization reflectometry of a hyperbolic metamaterial consisting of 8 layers of (Sn 5.3 nm)/(CaF₂ 1 nm). Data points are the measured p- and s-polarized reflectivities of the metamaterial sample at 488 nm, 633 nm and 1550 nm. The lines are fits using Eqs. (11-14). The fitting parameters, the real and imaginary parts of ε_1 and ε_2 , are specified in the plot. The negative sign of the real part of ε_1 and the positive sign of the real part of ε_2 obtained from the fits confirm hyperbolic character of the metamaterial.

Fig. 7. The calculated plots of the real and imaginary parts of (a) ε_1 and (b) ε_2 of the multilayer Sn/CaF₂ metamaterial. The calculations were performed using Eqs. (9, 10) based on the Maxwell-Garnett approximation. The metamaterial appears to be hyperbolic in the visible and near IR ranges. The experimentally measured data points from Fig. 6 are also shown for the comparison.

Fig. 8. Measured resistivity as a function of temperature for several multilayer Sn/CaF₂ metamaterials having different thicknesses of Sn layers, while the 1 nm thickness of CaF₂ layers was kept approximately constant. The thicknesses of Sn/CaF₂ layer are indicated. The Sn/CaF₂ hyperbolic metamaterial samples were found to have a higher critical temperature than pure Sn, which strongly indicates the role of the hyperbolic geometry in T_c enhancement.

Fig. 9. Effect of the tin volume fraction n on T_c of the Sn/CaF₂ and Sn/MgF₂ hyperbolic metamaterial samples. Experimentally measured behavior of T_c as a function of n (which is defined by the tin layer thickness) correlates well with the theoretical fit (red curve) based on the hyperbolic mechanism of T_c enhancement. Experimental data

points shown in black correspond to Sn/CaF₂ samples, while blue ones correspond to Sn/MgF₂ samples.

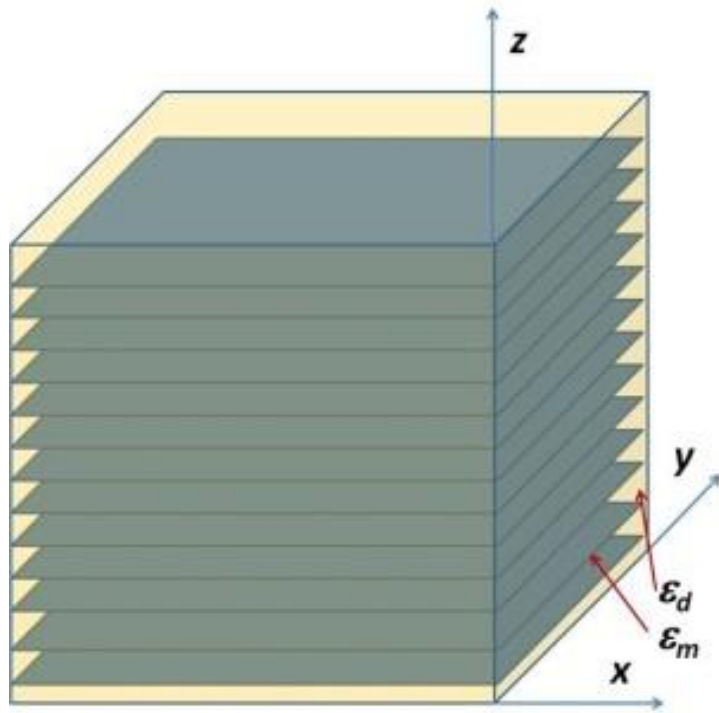
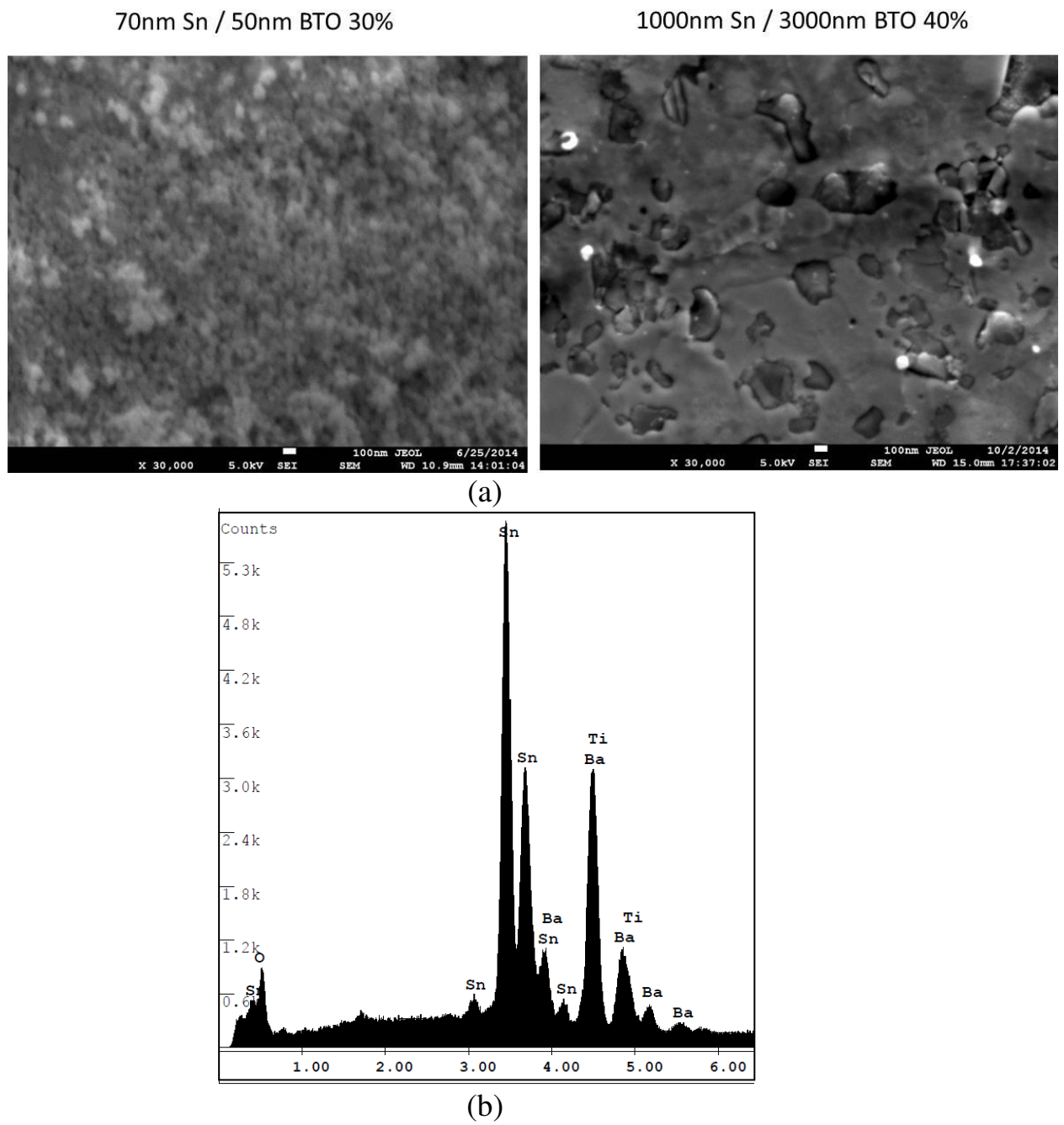
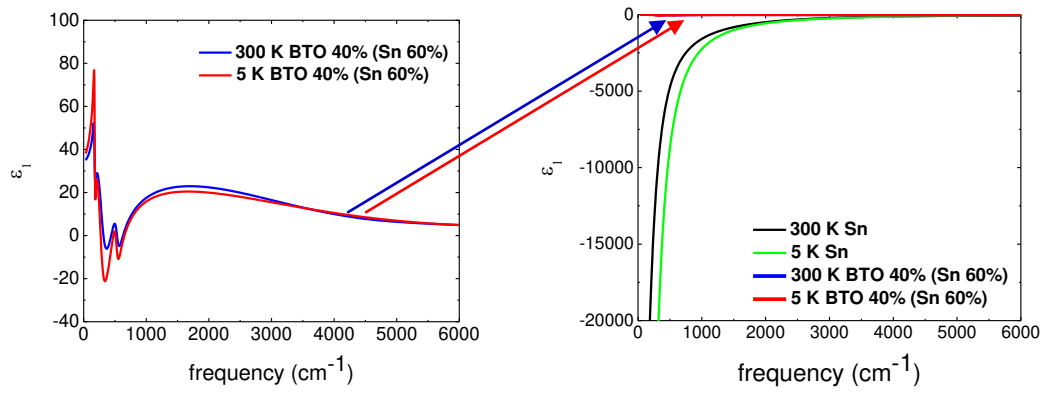
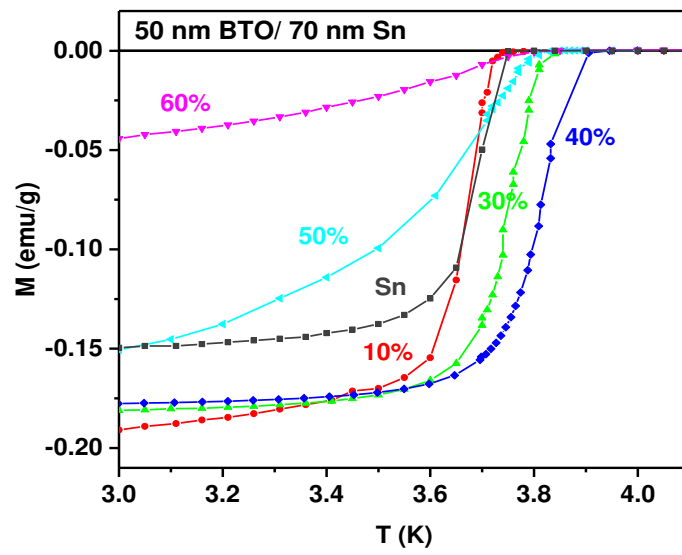


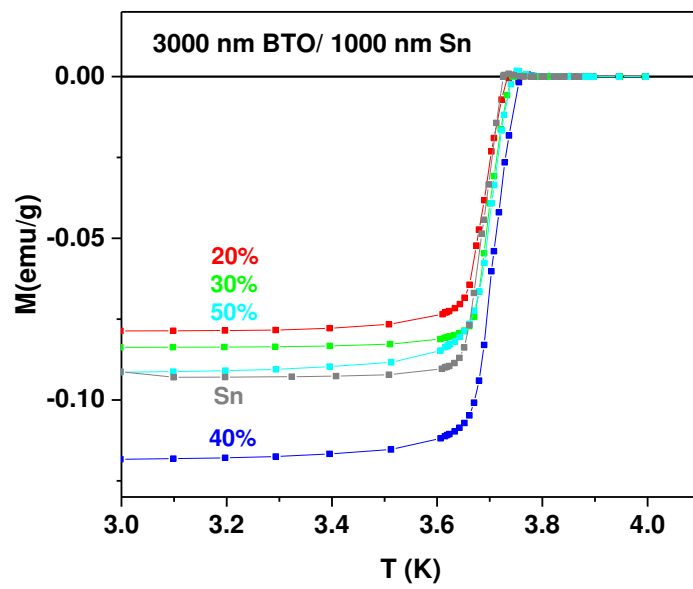
Fig. 1

**Fig. 2**

**Fig. 3**



(a)



(b)

Fig. 4

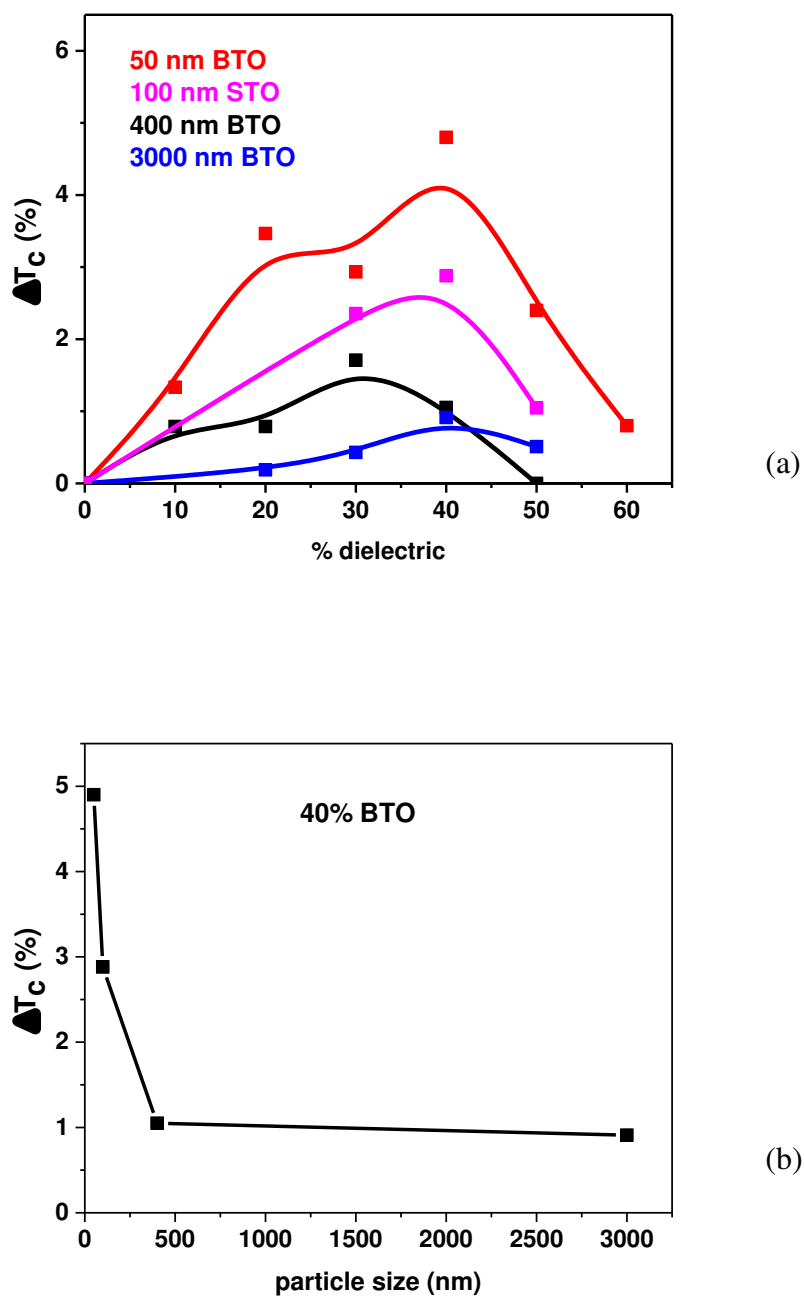


Fig. 5

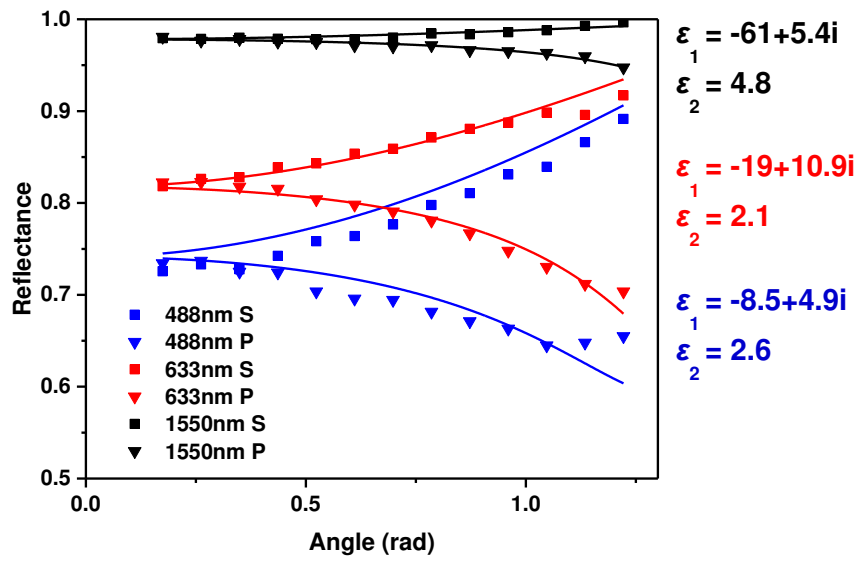


Fig. 6

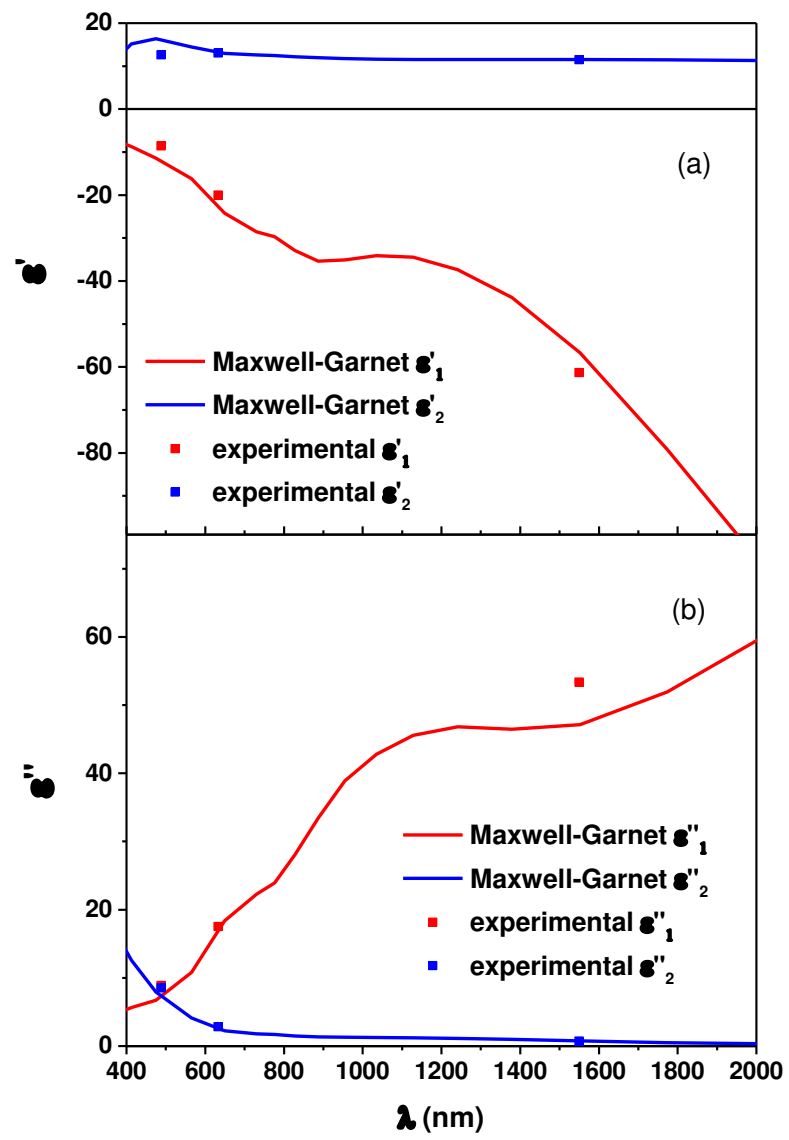


Fig. 7

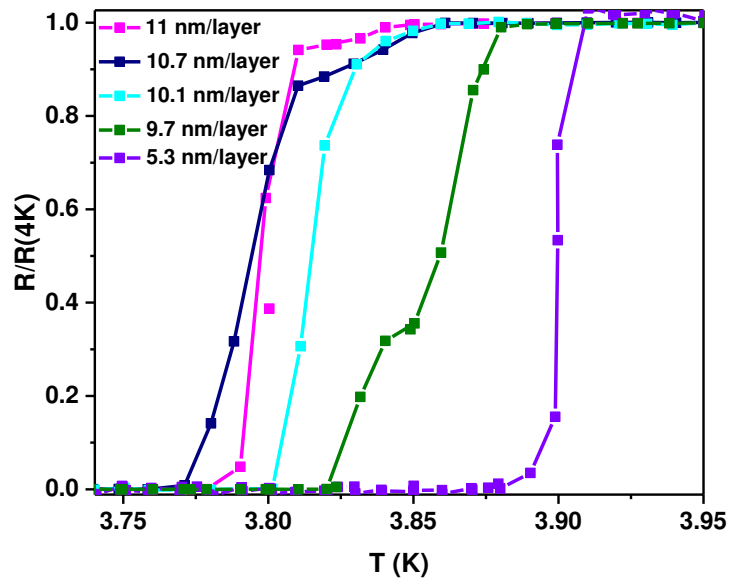
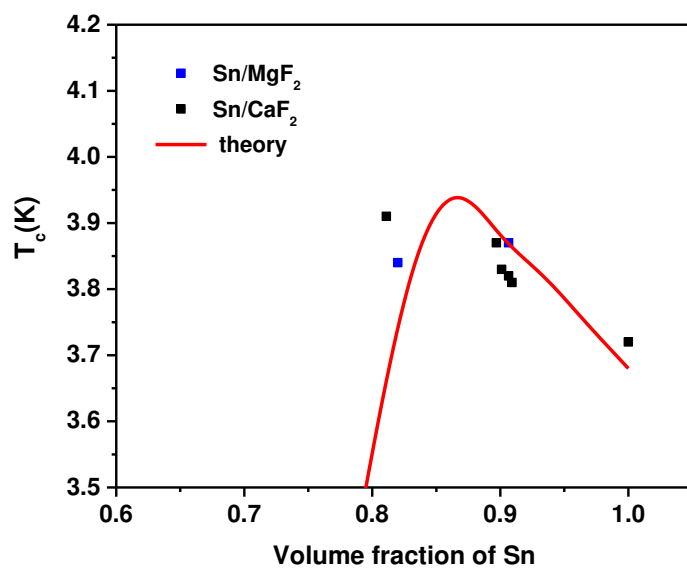


Fig. 8

**Fig. 9**